## Light bends!

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Solution to Maxwell's Equations

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www.physics.gatech.edu/frog/lectures

## Diffraction

Light does not always travel in a straight line.

It tends to bend around objects. This tendency is called diffraction.

Any wave will do this, including matter waves and acoustic waves.


Shadow of a zinc oxide crystal illuminated by a electrons


## Why it's hard to see diffraction

Diffraction tends to cause ripples at edges. But a point source is required to see this effect. A large source masks them.


Rays from a point source yield a perfect shadow of the hole. Rays from other regions blur the shadow.

Example: a large source (like the sun) casts blurry shadows, masking the diffraction ripples.

## Diffraction of ocean water waves

Ocean waves passing through slits in Tel Aviv, Israel


Diffraction occurs for all waves, whatever the phenomenon.

## Diffraction of a wave by a slit

Whether waves in water or electromagnetic radiation in air, passage through a slit yields a diffraction pattern that will appear more dramatic as the size of the slit approaches the wavelength of the wave.


## Diffraction by an Edge



Even without a small slit, diffraction
can be strong.

Light passing by edge
Simple propagation past an edge yields an unintuitive irradiance pattern.

Electrons passing by an edge (Mg0 crystal)


## Radio waves diffract around mountains.

When the wavelength is a km long, a mountain peak is a very sharp edge!

Another effect that occurs is scattering, so diffraction's role is not obvious.

DIFFRACTION


IROPOSCATIER


## Diffraction Geometry

We wish to find the light electric field after a screen with a hole in it.
This is a very general problem with far-reaching applications.


What is $E\left(x_{1}, y_{1}\right)$ at a distance $z$ from the plane of the aperture?

## Diffraction Assumptions

The best assumptions were determined by Kirchhoff:

1) Maxwell's equations

2) Inside the aperture, the field and its spatial derivative are the same as if the screen were not present.
3) Outside the aperture (in the shadow of the screen), the field and its spatial derivative are zero.

While these assumptions give the best results, they actually over-determine the problem and can be shown to yield zero field everywhere! Nevertheless, we still use them.

## Diffraction Solution

The field in the observation plane, $E\left(x_{1}, y_{1}\right)$, at a distance $z$ from the aperture plane is given by a convolution:

$$
E\left(x_{1}, y_{1}\right)=\iint h\left(x_{1}-x, y_{1}-y\right) t(x, y) E(x, y) d x d y
$$

where:

$$
h\left(x_{1}-x, y_{1}-y\right)=\frac{1}{i \lambda} \frac{\exp (i k r)}{r}
$$

$$
\longleftarrow \begin{aligned}
& \text { Spherical } \\
& \text { wave! }
\end{aligned}
$$

and:

$$
r=\sqrt{z^{2}+\left(x-x_{1}\right)^{2}+\left(y-y_{1}\right)^{2}}
$$

A very complicated result!

## Huygens' Principle

Huygens'
Principle says
that every
point along a wave-front emits a
spherical wave that interferes with all others.


Christiaan Huygens
1629-1695

Our solution for diffraction illustrates this idea, and it's more rigorous.

## Fresnel Diffraction: Approximations

In the denominator, we can approximate $r$ by $z$. But we can't approximate $r$ in the exp by $z$ because it gets multiplied by $k$, which is big, so relatively small changes in $r$ can make a big difference! But we can write:

$$
r=\sqrt{z^{2}+\left(x-x_{1}\right)^{2}+\left(y-y_{1}\right)^{2}}=z \sqrt{1+\left(\frac{x-x_{1}}{z}\right)^{2}+\left(\frac{y-y_{1}}{z}\right)^{2}}
$$

And if $\varepsilon \ll 1, \sqrt{1+\varepsilon} \approx 1+\varepsilon / 2$

$$
r \approx z\left[1+\frac{1}{2}\left(\frac{x-x_{1}}{z}\right)^{2}+\frac{1}{2}\left(\frac{y-y_{1}}{z}\right)^{2}\right]=z+\frac{\left(x-x_{1}\right)^{2}}{2 z}+\frac{\left(y-y_{1}\right)^{2}}{2 z}
$$

This yields: $\quad E\left(x_{1}, y_{1}\right)=$

$$
\iint \frac{1}{i \lambda z} \exp \left\{i k\left[z+\frac{\left(x-x_{1}\right)^{2}}{2 z}+\frac{\left(y-y_{1}\right)^{2}}{2 z}\right]\right\} t(x, y) E(x, y) d x d y
$$

## Fresnel Diffraction: Approximations

Multiplying out the squares: $\quad E\left(x_{1}, y_{1}\right)=$
$\iint \frac{1}{i \lambda z} \exp \left\{i k\left[z+\frac{\left(x^{2}-2 x x_{1}+x_{1}^{2}\right)}{2 z}+\frac{\left(y^{2}-2 y y_{1}+y_{1}^{2}\right)}{2 z}\right]\right\} t(x, y) E(x, y) d x d y$
Factoring out the quantities independent of $x$ and $y$ : $\quad E\left(x_{1}, y_{1}\right)=$

$$
\frac{\exp (i k z)}{i \lambda z} \exp \left[i k \frac{x_{1}^{2}+y_{1}^{2}}{2 z}\right] \iint \exp \left\{i k\left[\frac{\left(-2 x x_{1}-2 y y_{1}\right)}{2 z}+\frac{\left(x^{2}+y^{2}\right)}{2 z}\right]\right\} t(x, y) E(x, y) d x d y
$$

This is the Fresnel integral.
It yields the light wave field at the distance $z$ from the screen.

## Diffraction Conventions

We'll typically assume that a plane wave is incident on the aperture.

$$
E(x, y)=\text { constant }
$$

It still has an $\exp [i(\omega t-k z)]$, but it's constant with respect to $x$ and $y$.

$$
E\left(x_{1}, y_{1}\right) \underbrace{\frac{\exp (i k z)}{i \lambda z} \exp \left[i k \frac{x_{1}^{2}+y_{1}^{2}}{2 z}\right] \iint \exp \left\{i k\left[\frac{\left(-2 x x_{1}-2 y y_{1}\right)}{2 z}+\frac{\left(x^{2}+y^{2}\right)}{2 z}\right]\right\} t(x, y) d x d y .}_{~}
$$

And we'll usually ignore the various factors in front:

$$
E\left(x_{1}, y_{1}\right) \propto \iint_{-\infty} \int_{-\infty} \exp \left\{i k\left[\frac{\left(-2 x x_{1}-2 y y_{1}\right)}{2 z}+\frac{\left(x^{2}+y^{2}\right)}{2 z}\right]\right\} t(x, y) d x d y
$$

## Fresnel diffraction: example

Fresnel diffraction from a single slit:


## Fresnel Diffraction from a Slit

This irradiance vs. position just after a slit illuminated by a laser.


## The Spot of Arago

If a beam encounters a stop, it develops a hole, which fills in as it propagates and diffracts:


This irradiance can be quite high and can do some damage!

## Fresnel diffraction from an array of slits: The Talbot Effect

One of the few Fresnel diffraction problems that can be solved analytically is an array of slits.
The beam pattern alternates between two different fringe patterns.


## The Talbot Carpet

What goes on in between the solvable planes?

The beam propagates in this direction.

The slits are here.


## Diffraction Approximated

These integrals come up:

$$
\begin{aligned}
C(x) & =\int_{0}^{x} \cos \left(t^{2}\right) d t \\
S(x) & =\int_{0}^{x} \sin \left(t^{2}\right) d t
\end{aligned}
$$

Such effects can be modeled by measuring the distance on a Cornu Spiral.

But most useful diffraction effects do not occur in the Fresnel diffraction regime because it's too complex.
For a cool Java applet that computes Fresnel diffraction patterns, try http://falstad.com/diffraction/

## Fraunhofer Diffraction: The Far Field

Recall the Fresnel diffraction result: $\quad E\left(x_{1}, y_{1}\right)=$

$$
\frac{\exp (i k z)}{i \lambda z} \exp \left[i k \frac{x_{1}^{2}+y_{1}^{2}}{2 z}\right] \iint \exp \left\{i k\left[\frac{\left(-2 x x_{1}-2 y y_{1}\right)}{2 z}+\frac{\left(x^{2}+y^{2}\right)}{2 z}\right]\right\} t(x, y) E(x, y) d x d y
$$

Let $D$ be the size of the aperture: $D^{2} \geq x^{2}+y^{2}$.
When $k D^{2} / 2 z \ll 1$, the quadratic terms $\ll 1$, so we can neglect them:

$$
E\left(x_{1}, y_{1}\right)=\frac{\exp (i k z)}{i \lambda z} \exp \left[i k \frac{x_{1}^{2}+y_{1}^{2}}{2 z}\right] \iint \exp \left\{-\frac{i k}{z}\left(x x_{1}+y y_{1}\right)\right\} t(x, y) E(x, y) d x d y
$$

This condition means going a distance away: $z \gg k D^{2} / 2=\pi D^{2} / \lambda$ If $D=1 \mathrm{~mm}$ and $\lambda=1$ micron, then $z \gg 3 \mathrm{~m}$.

## Fraunhofer Diffraction Conventions

Neglect the phase factors, and we'll explicitly write the aperture transmission function, $t(x, y)$, in the integral:

$$
E\left(x_{1}, y_{1}\right) \propto \int_{-\infty}^{\infty} \exp \left\{-\frac{i k}{z}\left(x x_{1}+y y_{1}\right)\right\} t(x, y) E(x, y) d x d y
$$

## This is just a Fourier Transform!

$$
E(x, y)=\text { constant if a plane wave }
$$

Interestingly, it's a Fourier Transform from position, $x$, to another position variable, $x_{1}$ (in another plane). Usually, the Fourier "conjugate variables" have reciprocal units (e.g., $t \& \omega$, or $x \& k$ ). The conjugate variables here are really $x$ and $k_{x}=k x_{1} / z$, which have reciprocal units.

So the far-field light field is the Fourier Transform of the transmitted field!

## The Fraunhofer Diffraction formula

We can write this result in terms of the off-axis $k$-vector components:

$$
E\left(k_{x}, k_{y}\right) \propto \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \exp \left[-i\left(k_{x} x+k_{y} y\right)\right] t(x, y) E(x, y) d x d y
$$

that is:

$$
E\left(k_{x}, k_{y}\right) \propto \mathscr{F}\{t(x, y) E(x, y)\}
$$

and:

$$
k_{x}=k x_{1} / z \text { and } k_{y}=k y_{1} / z
$$

or:

$$
\theta_{x}=k_{x} / k=x_{1} / z \text { and } \theta_{y}=k_{y} / k=y_{1} / z
$$



## The Uncertainty Principle in Diffraction!

$$
E\left(k_{x}, k_{y}\right) \propto \mathscr{F}\{t(x, y) E(x, y)\} \quad k_{x}=k x_{1} / z
$$

Because the diffraction pattern is the Fourier transform of the slit, there's an uncertainty principle between the slit width and diffraction pattern width!

If the input field is a plane wave and $\Delta x$ is the slit width and $\Delta k_{x}$ is the proportional to the beam angular width after the screen,

$$
\Delta x \Delta k_{x}>1
$$

Or:

$$
\Delta x \Delta x_{1}>z / k
$$

The smaller the slit, the larger the diffraction angle and the bigger the diffraction pattern!

## Fraunhofer Diffraction from a slit

Fraunhofer Diffraction from a slit is simply the Fourier Transform of a rect function, which is a sinc function. The irradiance is then sinc ${ }^{2}$.

$$
\begin{aligned}
& t(x)=\operatorname{rect}[x / w] \\
& E\left(k_{x}\right) \propto \mathscr{F}\{t(x)\}
\end{aligned}
$$

$$
\begin{aligned}
& E\left(k_{x}\right) \propto \operatorname{sinc}\left(w k_{x} / 2\right) \\
& E\left(x_{1}\right) \propto \operatorname{sinc}\left(w k x_{1} / 2 z\right)
\end{aligned}
$$

$$
\begin{aligned}
& I\left(k_{x}\right) \propto \operatorname{sinc}^{2}\left(w k_{x} / 2\right) \\
& I\left(x_{1}\right) \propto \operatorname{sinc}^{2}\left(w k x_{1} / 2 z\right)
\end{aligned}
$$

## Fraunhofer Diffraction from a Square Aperture

The diffracted field is a sinc function in both $x_{1}$ and $y_{1}$ because the Fourier transform of a rect function is sinc.



Diffracted irradiance


Diffracted field

## Diffraction from a Circular Aperture

A circular aperture yields a diffracted "Airy Pattern," which looks a lot like a sinc function, but actually involves a Bessel function.


Diffracted field


Diffracted Irradiance

## Diffraction from

small and large circular apertures

Far-field intensity pattern from a small aperture



Recall the Scale Theorem! This is the Uncertainty Principle for diffraction.



$t(x)=\operatorname{rect}[(x+a) / w]+\operatorname{rect}[(x-a) / w]$

$$
E\left(k_{x}\right) \propto \mathscr{F}\{t(x)\}
$$

$\propto \operatorname{sinc}\left[w\left(k x_{1} / z\right) / 2\right] \exp \left[+i a\left(k x_{1} / z\right)\right]+$ $\operatorname{sinc}\left[w\left(k x_{1} / z\right) / 2\right] \exp \left[-i a\left(k x_{1} / z\right)\right]$
$E\left(x_{1}\right) \propto \operatorname{sinc}\left(w k x_{1} / 2 z\right) \cos \left(a k x_{1} / z\right)$


## Diffraction from one- and two-slit screens

Fraunhofer diffraction patterns

One slit

Two slits


## Diffraction from multiple slits

Infinitely many equally spaced slits (a Shah function!) yields a far-field pattern that's the Fourier transform, that is, the Shah function.


## Two Slits and Spatial Coherence



If the spatial coherence length is less than the slit separation, then the relative phase of the light transmitted through each slit will vary randomly, washing out the fine-scale fringes, and a one-slit pattern will be observed.

Fraunhofer diffraction patterns


## Young's Two Slit Experiment and Quantum Mechanics

Imagine using a beam so weak that only one photon passes through the screen at a time. In this case, the photon would seem to pass through only one slit at a time, yielding a one-slit pattern. Which pattern occurs?

Possible Fraunhofer diffraction patterns


## Dimming the light incident on two slits

Dimming the light in a two-slit experiment yields single photons at the screen. Since photons are particles, it would seem that each can only go through one slit, so then their pattern should become the single-slit pattern.


Each individual photon goes through both slits!

$$
\text { (a) } 20 \text { counts }
$$

